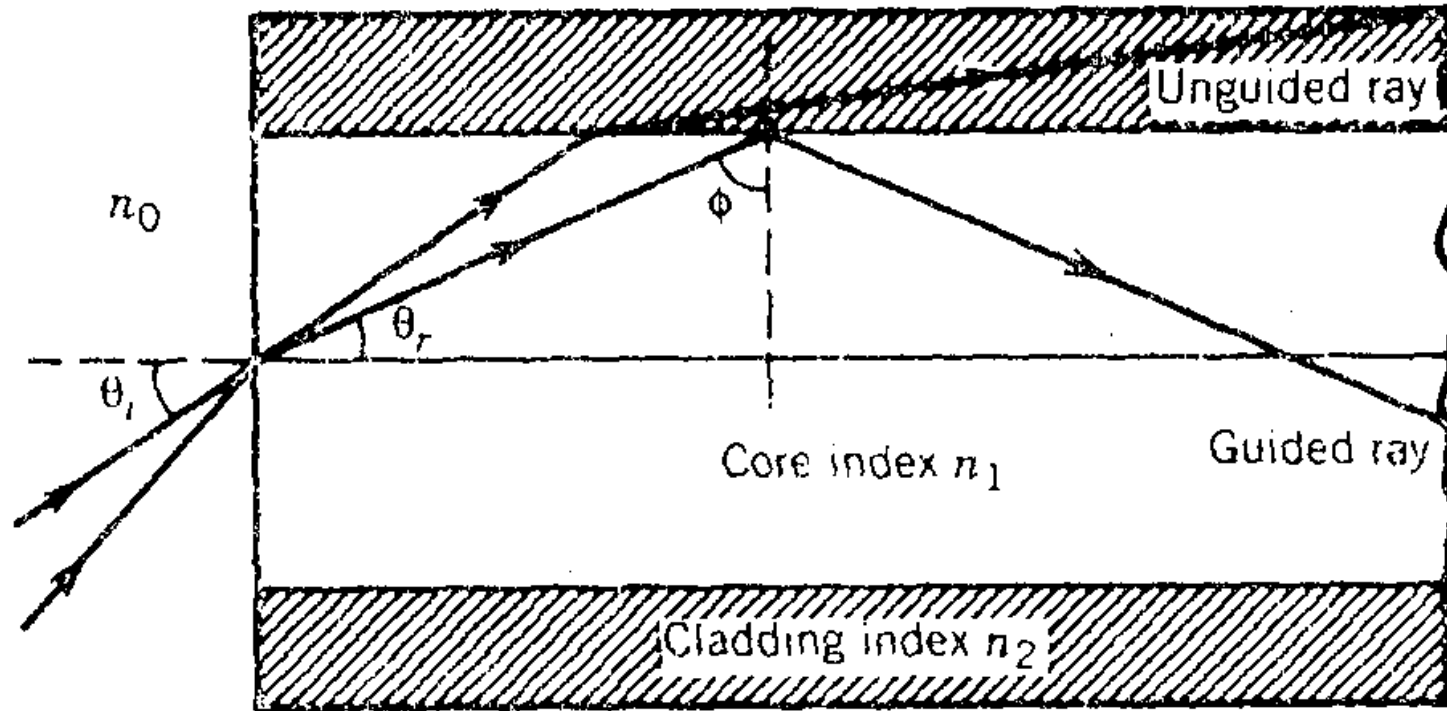




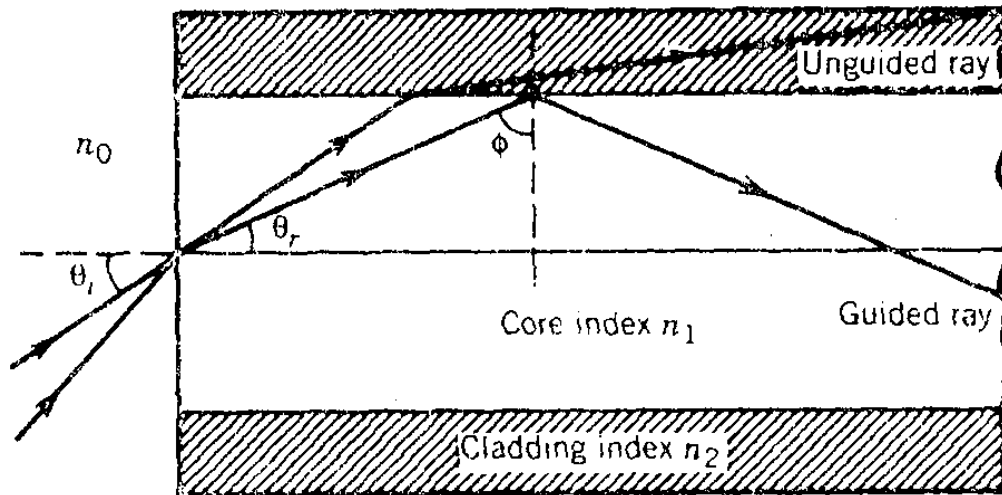
Slab Waveguides

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**Primary Reference: C. Pollock, *Fundamentals of Optoelectronics*,
Chapter 2, Richard Irwin Inc, (1995), ISBN 0-256-10104-3**



What is the maximum incident angle θ_1 that will enable waveguiding of the optical ray?



- $\theta_{1,\max}$ is a function of $\phi_{1,\min}$. Why?
- $\phi_{1,\min}$ is actually ϕ_{TIR} !
- $n_1 \sin(\phi_{\text{TIR}}) = n_2 \sin(\pi/2)$
- For coupling into the waveguide, $\sin\theta_i$ has to be less than $n_1 \sin\theta_r$ (via Snell' s Law)
- Therefore $\sin\theta_i < \sin\theta_{i,\max} = n_1 \sin(\pi/2 - \theta_{\text{TIR}})$
- Numerical Aperture = N.A. $\equiv \sin\theta_{1,\max}$



- Numerical Aperture = N.A. $\equiv \sin\theta_{i,\max} = n_1 \sin(\pi/2 - \phi_{\text{TIR}}) = n_1 \cos \phi_{\text{TIR}}$
- $n_1 \sin(\phi_{\text{TIR}}) = n_2 \sin(\pi/2) = n_2$
- $[n_1 \sin(\phi_{\text{TIR}})]^2 = [n_2]^2 \rightarrow \sin^2(\phi_{\text{TIR}}) = [n_2/n_1]^2$
- $\sin^2(\phi_{\text{TIR}}) + \cos^2(\phi_{\text{TIR}}) = 1 \rightarrow [n_2/n_1]^2 + \cos^2(\phi_{\text{TIR}}) = 1$

Now we are ready to develop an expression for NA ...

$$\mathbf{N.A.} = n_1 \cos(\theta_{\text{TIR}}) = n_1 (1 - [n_2/n_1]^2)^{1/2} = \mathbf{(n_1^2 - n_2^2)^{1/2}}$$

$$\Delta = \frac{n_1 - n_2}{n_1} \quad \text{Let's assume the case when } n_1 \approx n_2 \text{ s.t. } \Delta \ll 1$$

$$(n_1 - n_2)^{1/2} (n_1 + n_2)^{1/2} = (\Delta n_1 2 n_1)^{1/2} = n_1 (2 \Delta)^{1/2}$$

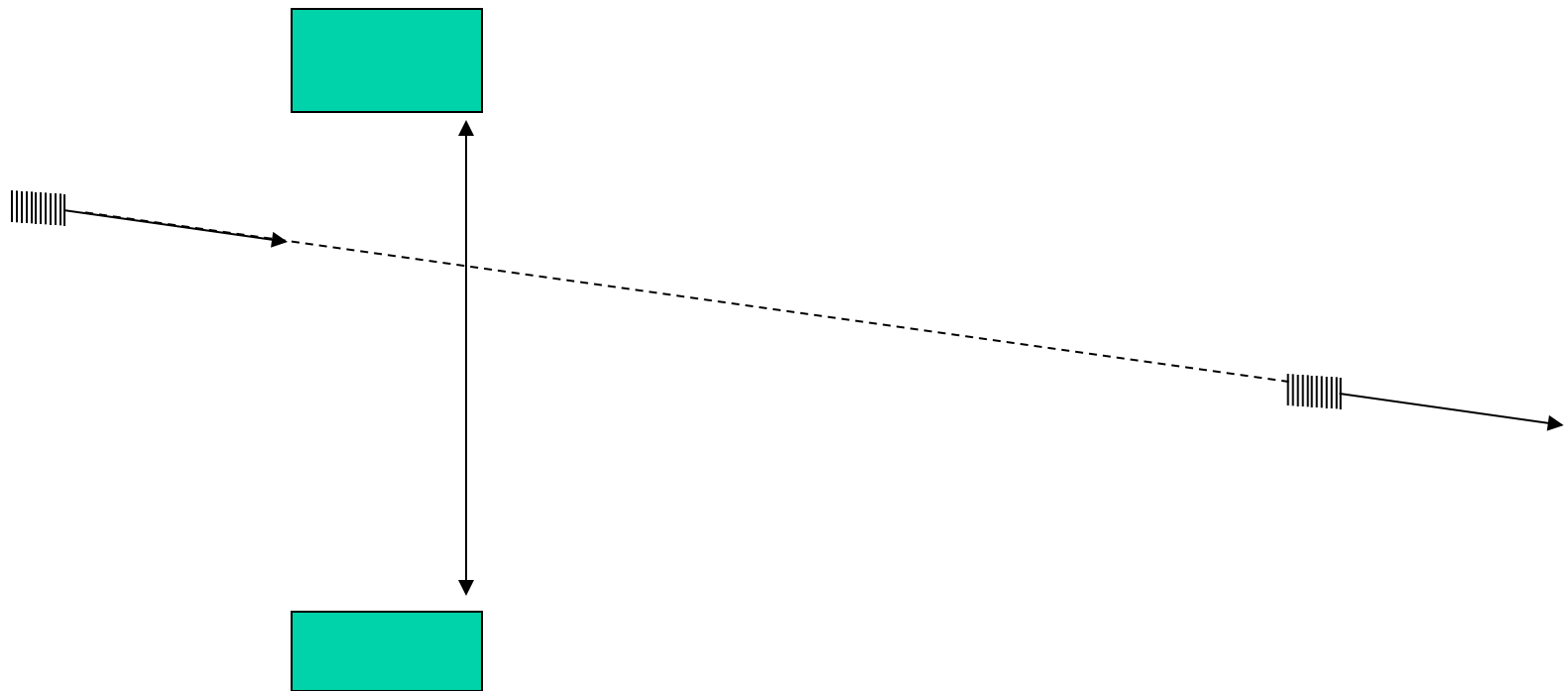
$$\text{Hence } \mathbf{N.A.} \approx \mathbf{n_1 (2 \Delta)^{1/2}}$$



n_1	n_2	Δ	NA	$\theta_i = \theta_{NA}$
1.446	1.440	4.149×10^{-3}	0.1316	7.56°
1.446	1.389	0.0394	0.40	23.6°
1.446	1.000	0.3084	1.044	?????



- Ray picture is inadequate to explain all observed phenomena in light propagation in waveguides
- The wave picture becomes significant when the waveguide dimensions approach the wavelength of light



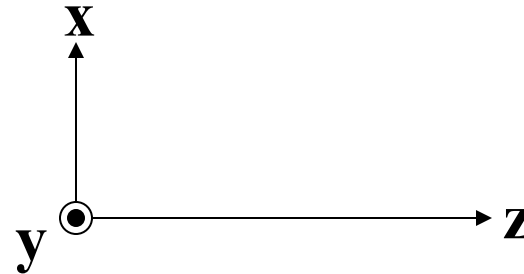


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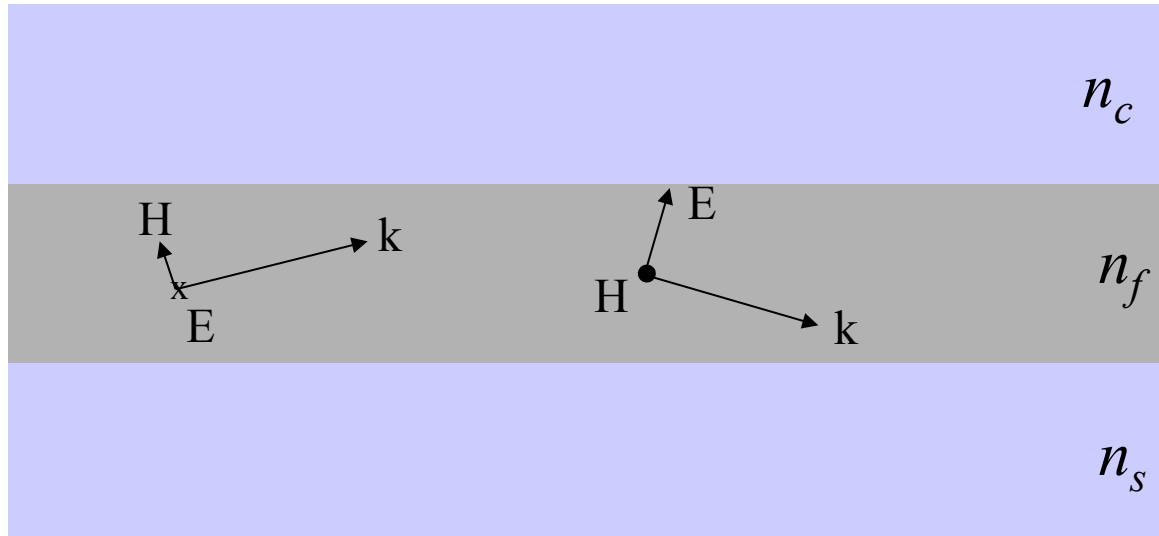




- Let's look at the planar slab waveguide
- Step-index/confinement in the x-direction
- Infinite in extent in the z and y directions



- Solve the wave equation with B.C.s at the interfaces
- This will lead to the concept of modes

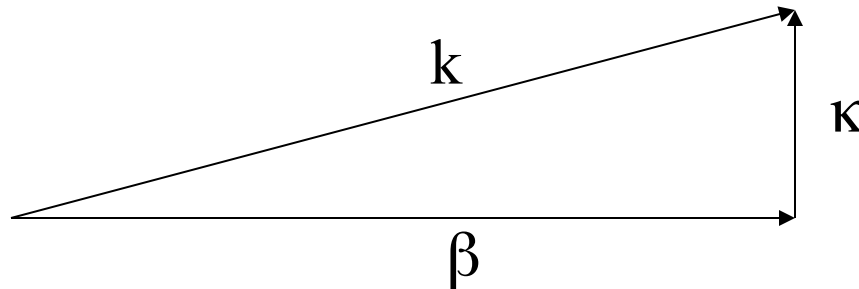


- +z is the direction of propagation
- Transverse Electric - no longitudinal E
- Transverse Magnetic - no longitudinal H
- $n_f > n_s > n_c$
- **TE case:** $\nabla^2 E_y + k_o^2 n_i^2 E_y = 0$, $n_i =$ function of layer, x, only



- **TE case:** $\nabla^2 E_y + k_o^2 n_i^2 E_y = 0$, n_i = function of layer, x , only
- $E_y(x,z) = E_y(x) \exp(-j\beta_i z)$, trial solution, beta is prop. coeff.
- Plugging into the wave equation, with $E_y'' = 0$ we get:

$$\nabla_{xx} E_y + (k_o^2 n_i^2 - \beta^2) E_y = 0$$
, n_i is a function of layer
- Decaying solution $k_o n_i < \beta$; $E_y(x) = E_o e^{-\gamma x}$; attenuation coeff.
- Oscillatory solution $k_o n_i > \beta$; $E_y(x) = E_o e^{\pm jkx}$; transverse wavevector
- Total wavevector k is the vector sum of the transverse wavevector and the propagation constant (longitudinal component)

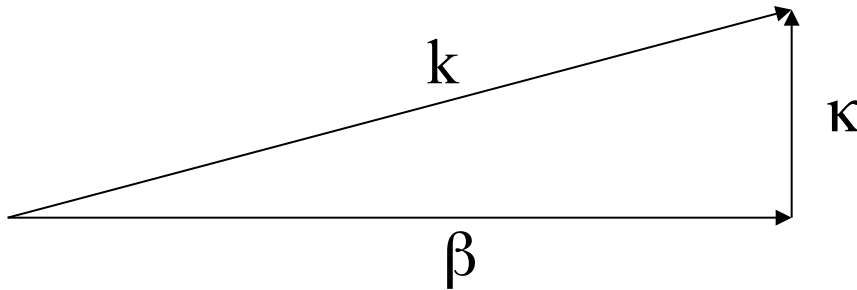


Wavevectors: $\nabla_{xx} \mathbf{E}_y + (k_o^2 n_i^2 - \beta^2) \mathbf{E}_y = 0$

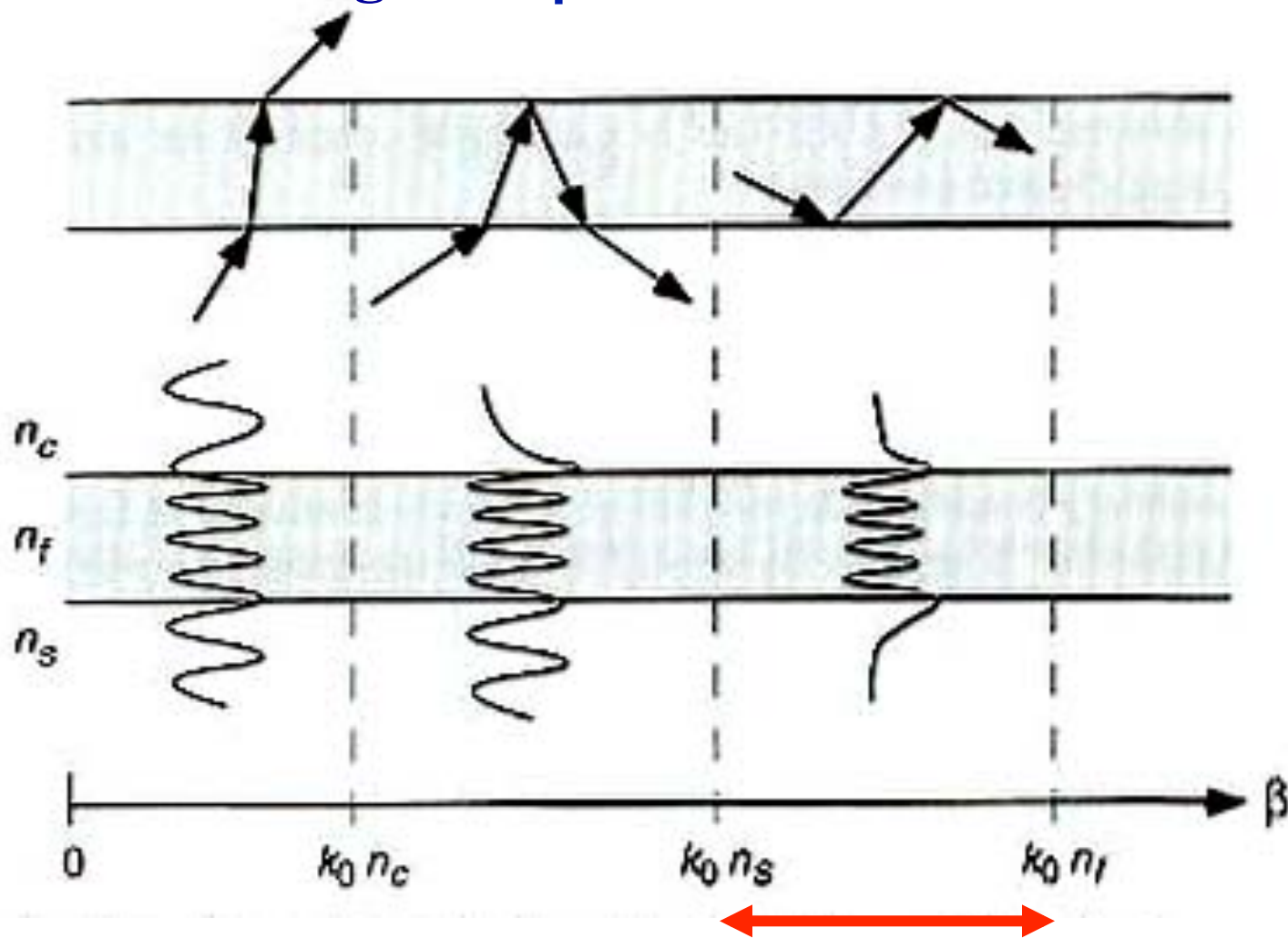


$$E_o e^{\pm j\kappa x} \implies \kappa = (k_o^2 n_i^2 - \beta^2)^{1/2}$$

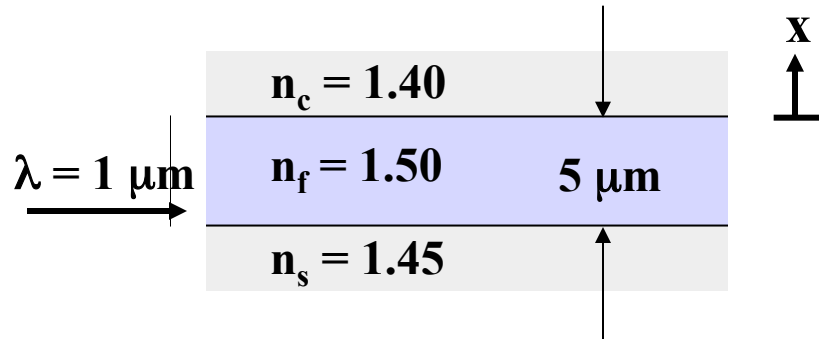
$$E_o e^{-\gamma x} \implies \gamma = (\beta^2 - k_o^2 n_i^2)^{1/2}$$



- wavevector
- transverse wavevector
- propagation constant

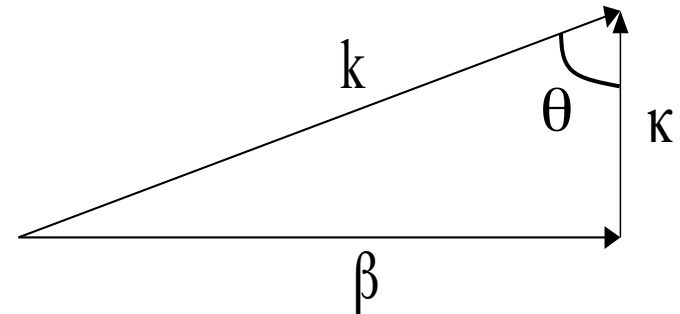


Varying the angle between k and z from 90 deg to 0 deg varies β

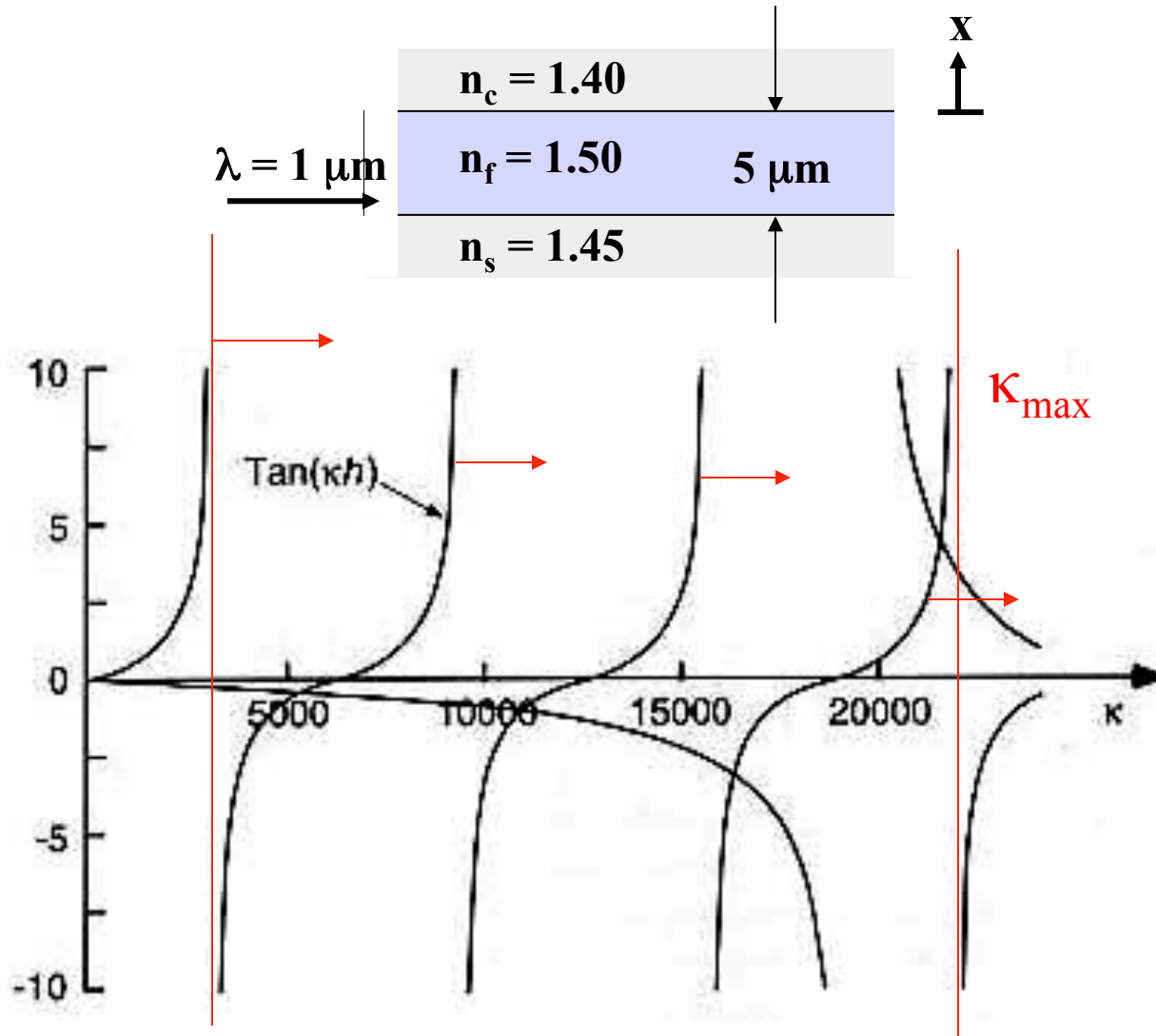


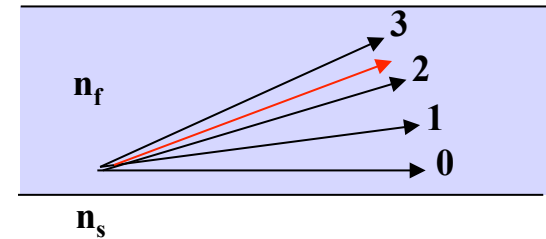
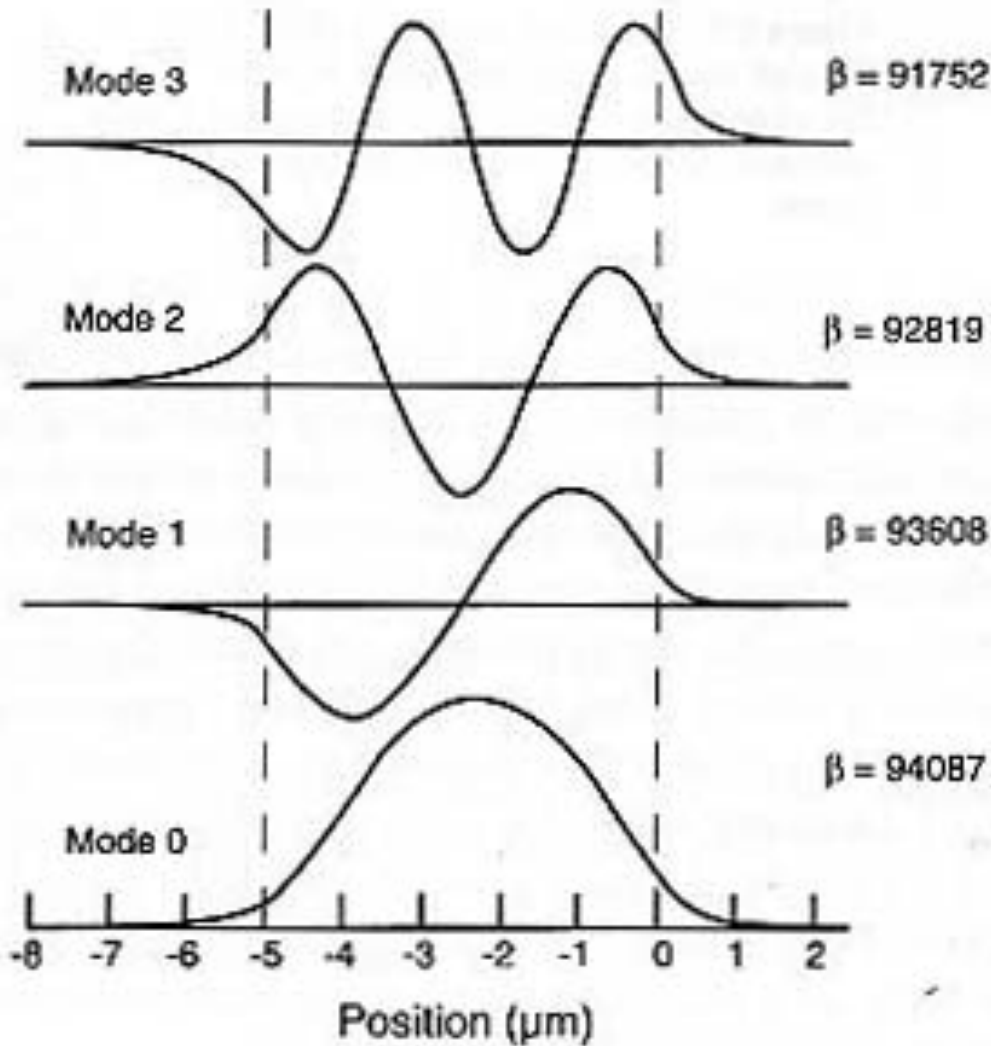
- $\gamma_c = \gamma_c(\kappa_f) = (\beta^2 - k_o^2 1.4^2)^{1/2}$
- $\gamma_s = \gamma_s(\kappa_f) = (\beta^2 - k_o^2 1.45^2)^{1/2}$
- $\beta = \beta(\kappa_f) = (k_o^2 1.5^2 - \kappa_f^2)^{1/2}$

- The minimum value of κ_f is 0, when the wavevector $k = k_o n_f = \beta$
- The maximum value of κ_f is $(k_o^2 n_f^2 - k_o^2 n_s^2)^{1/2}$



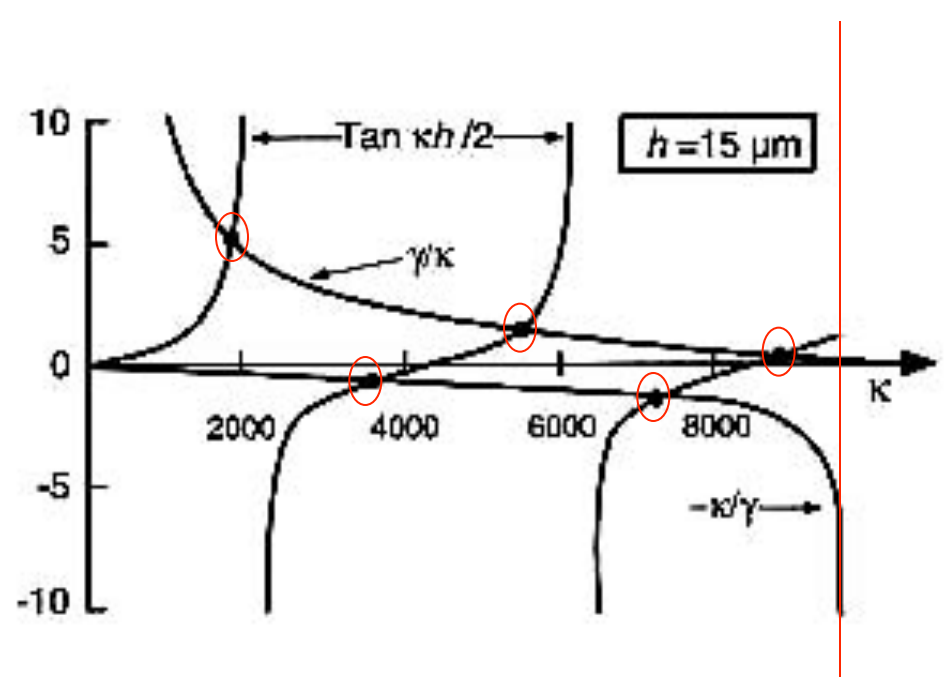
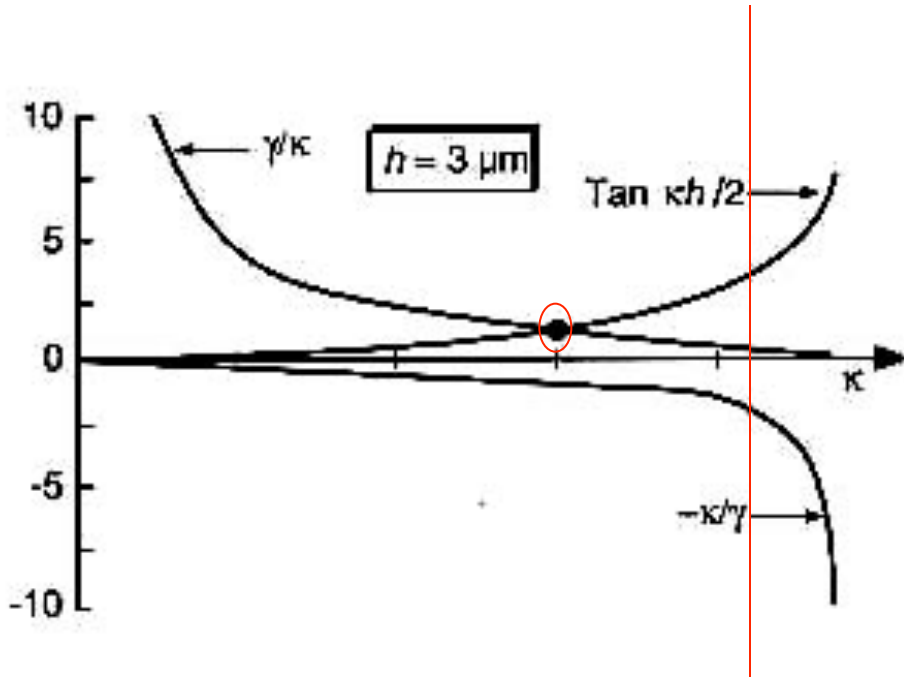
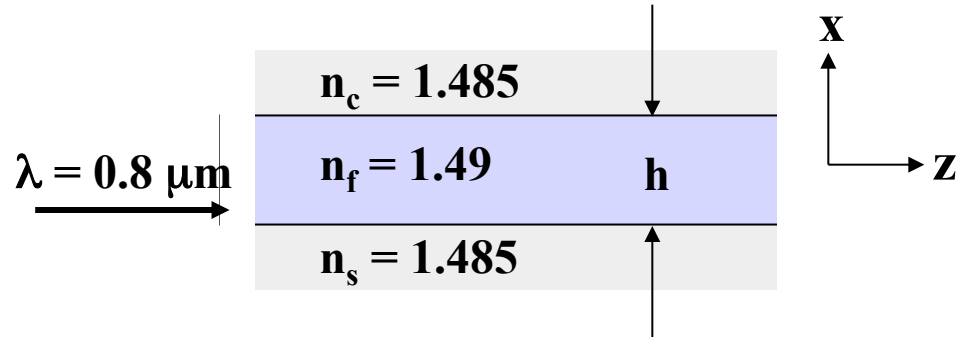
Slab WG Example

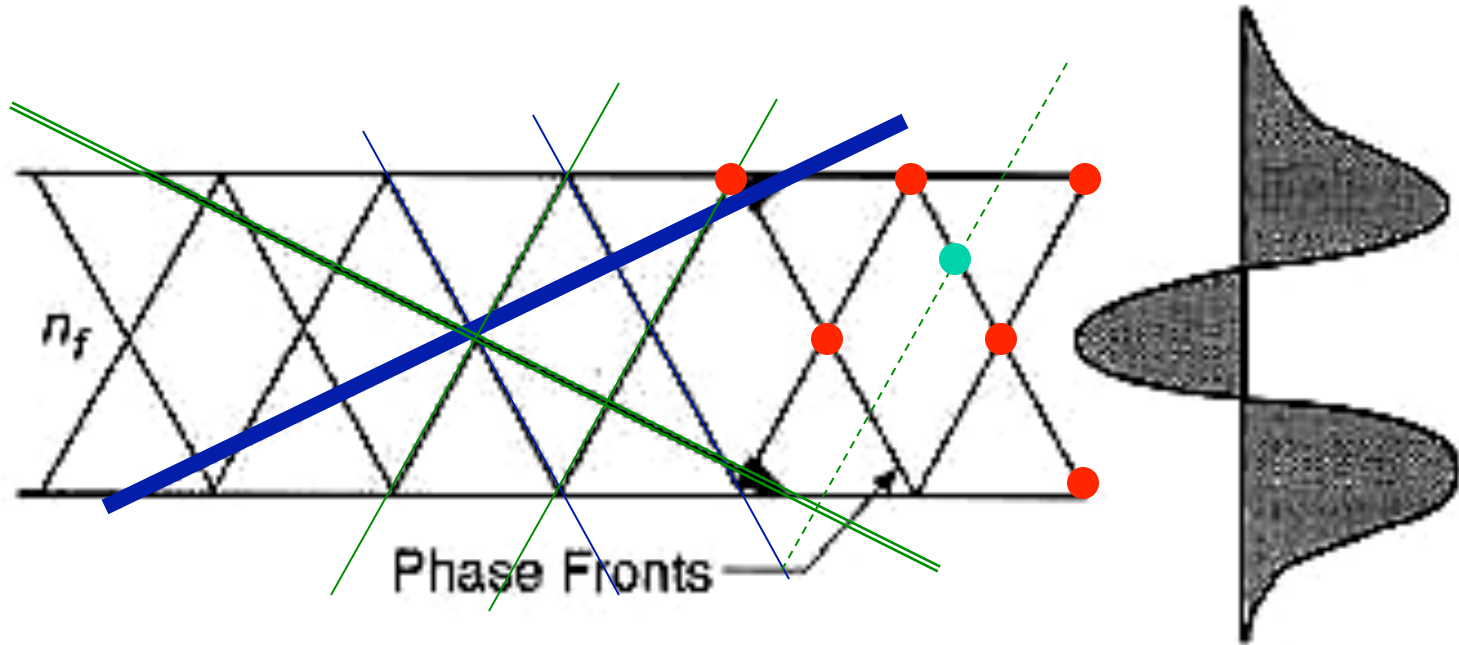




P_{core} / P_{total}	Mode
85.9 %	3
	2
	1
99.47 %	0

Symmetric Waveguide Example





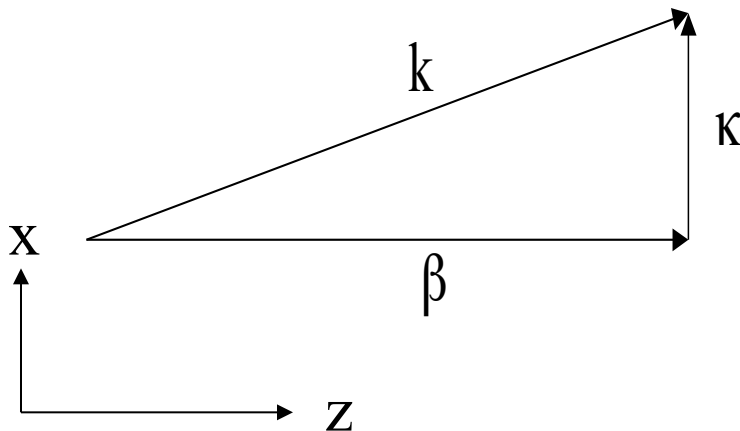


- The maximum value of β is $k_0 n_f$
- The minimum value of β is $k_0 n_s$
- There are only a *discrete number* of β in this range

This corresponds to a ray parallel to z into the waveguide

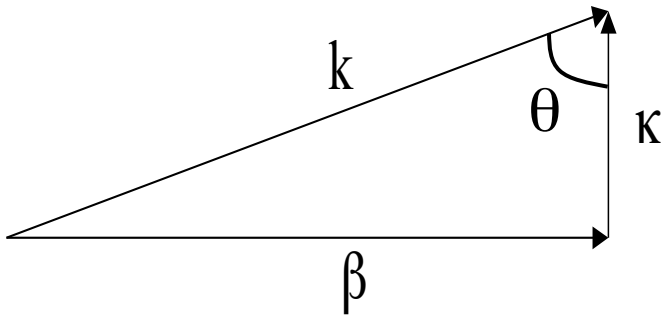
This corresponds to a ray at the critical angle to the film-substrate interface

A resonance condition must be satisfied if a mode is to be supported in the waveguide, only a discrete number of modes can satisfy this resonance condition





- *Confinement* in one dimension leads to
- *Standing waves* and confined *modes* in that dimension

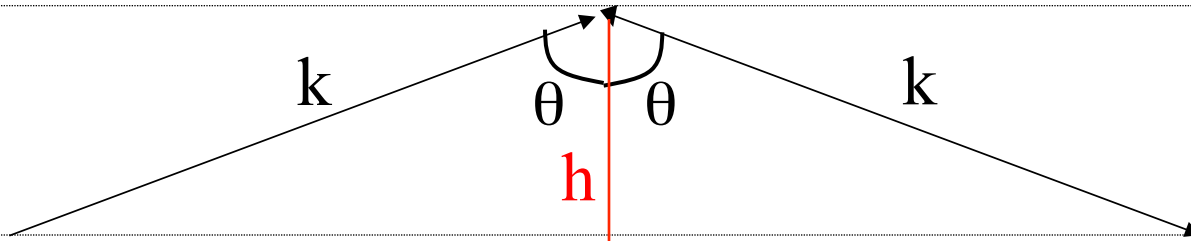


Phase shift = $k \cdot d = 2 \pi \nu$,
 ν is an integer

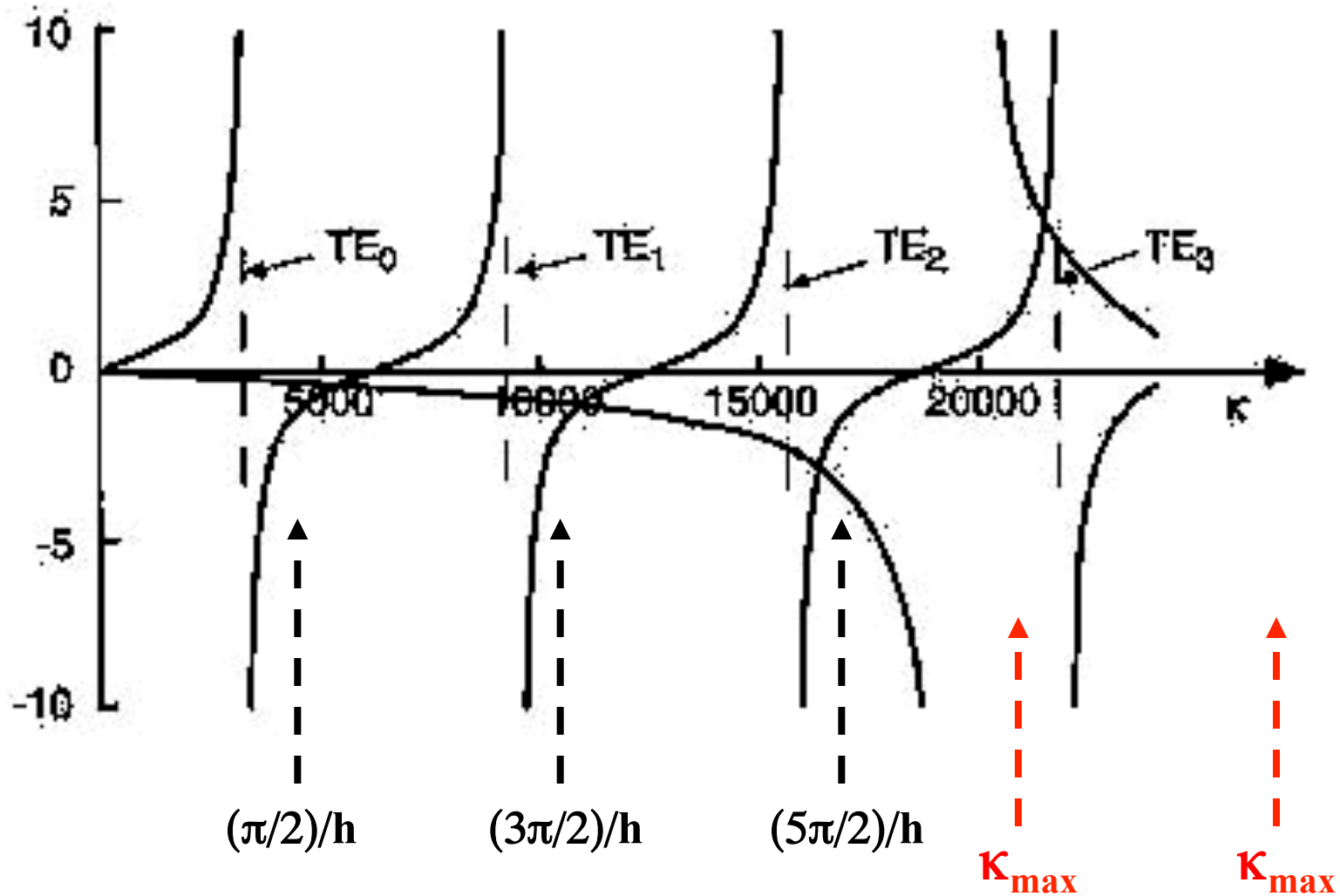
cover

film

substrate



$$2\pi\nu = 2 k_o n_f h \cos\theta - 2\Phi_c - 2\Phi_s$$



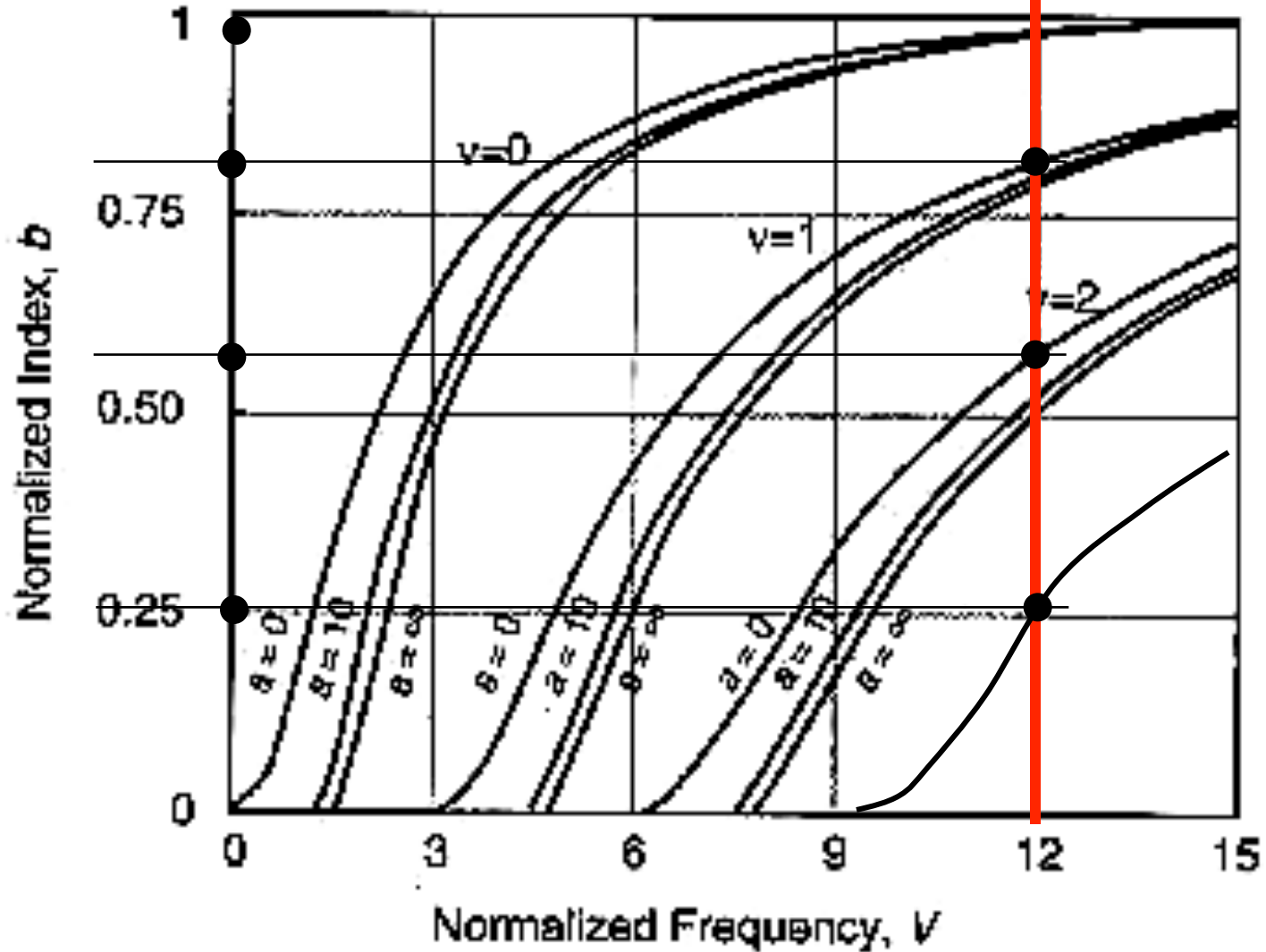
Pulse distortion and bandwidth are dependent on number of modes



- Approximation below is best applied when there is a large number of modes:

$$m = \text{Int}[\kappa_{\max} h / \pi] = \text{Int}[h k_o (n_f^2 - n_s^2)^{1/2} / \pi]$$

- Number of modes increase with thickness h , with index difference between film and cladding, and as light wavelength λ_o gets shorter.



Normalized Index: Slab Waveguide

